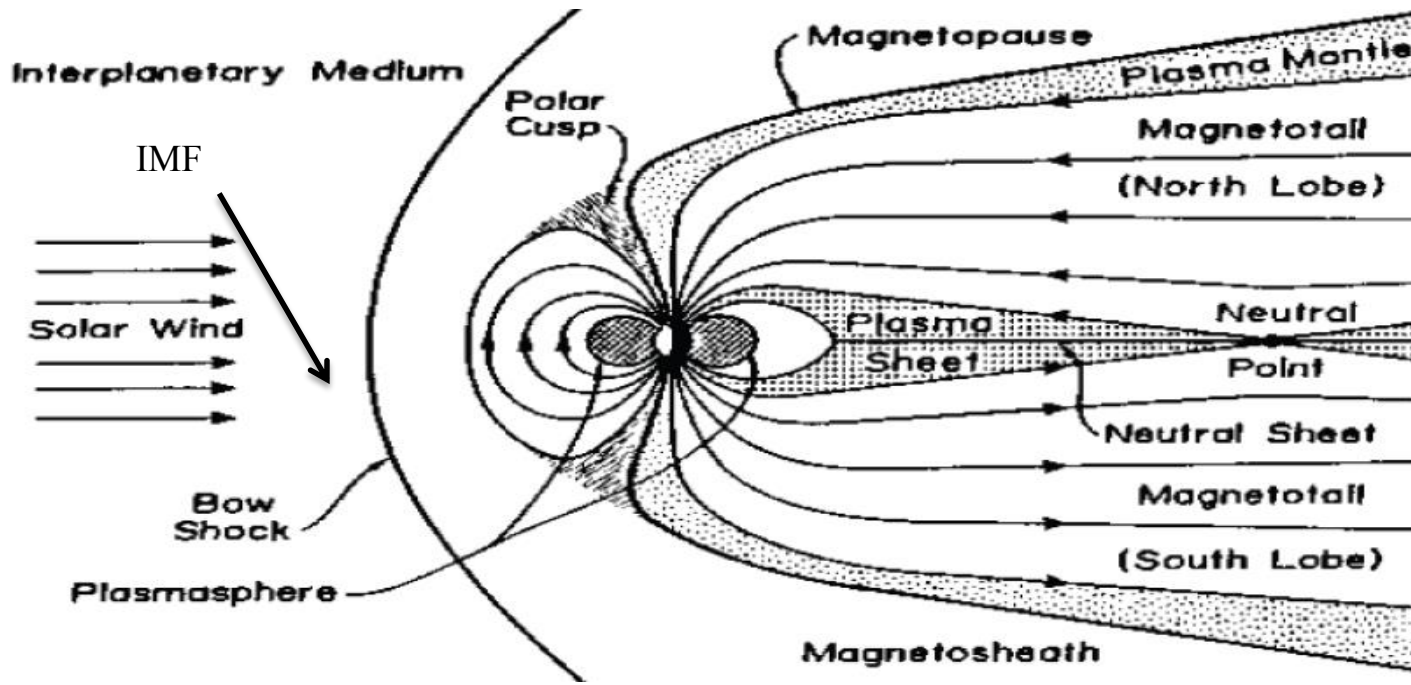


*International Space Sciences School*  
*Heliospheric physical processes for understanding Solar*  
*Terrestrial Relations*  
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**Lecture 1: Introduction to Space Plasma Physics**

A schematic diagram of the solar terrestrial plasma environment.



Understanding this diagram requires knowledge about

- transport of solar magnetic field
- how SW interacts with planetary magnetic fields
- collisionless shocks form
- the importance of neutral points
- current sheets
- particles acceleration mechanisms in the Magnetosphere.

## *Goal of the Lectures:*

- To help understand space plasma behavior, *we will provide useful material* not normally found in space plasmas textbooks (Russell and Kivelson, 1996; Parks, 1996; 2004) .
- Space plasma features are complex and often can have more than one interpretation.
- *Identify Issues with some models* and suggest different ways to interpret the data or how to resolve the issues.
- You *may not agree* with the ideas and concepts given in these lectures. *Criticisms, comments and questions* are welcome!

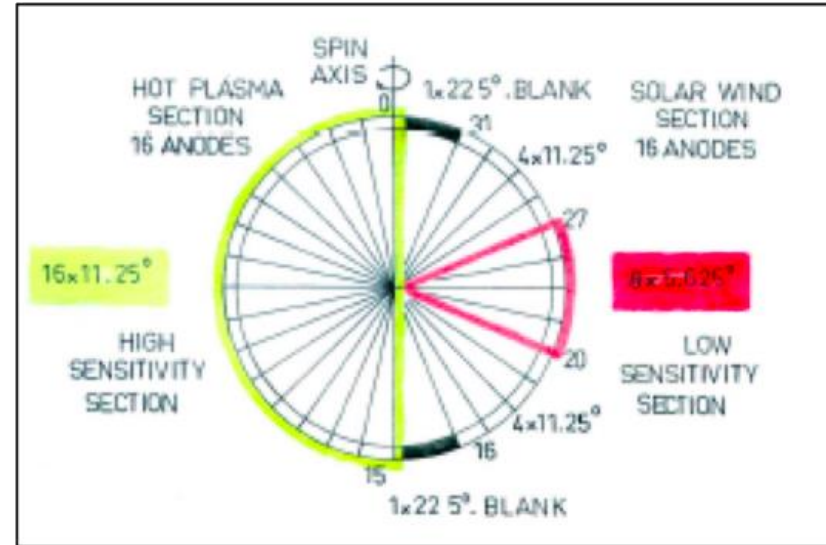
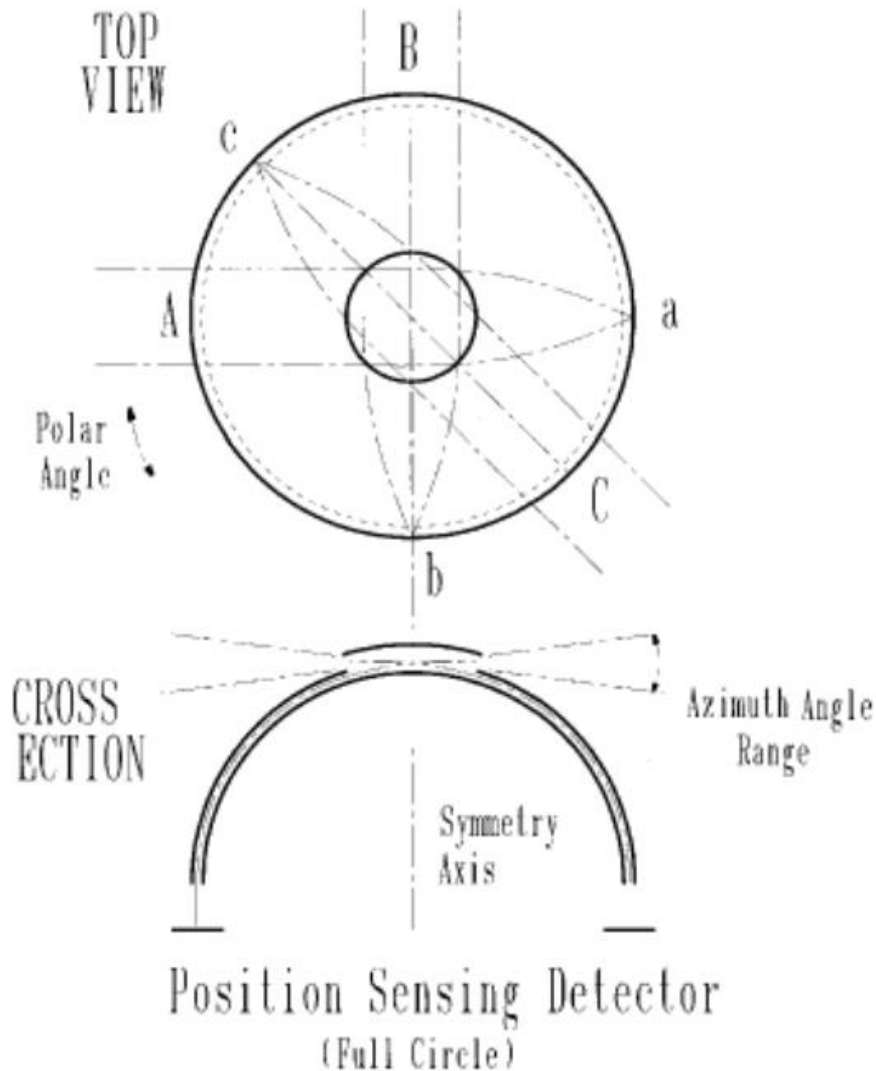
## *Point of View:*

- Information about space plasmas comes from measurements made by in-situ experiments. If there are disagreement about interpretation, we go back to data.
- This first lecture will briefly review
  - (1) how detectors work and what they measure,
  - (2) what assumptions are made in the measurements, which affect interpretation of data,
  - (3) basic plasma theories and concepts needed to interpret the data.

## Plasma Instrument and Measurements

- Focus on space plasmas with energies a few eV to  $\sim 40$  keV/charge which includes most of solar wind and magnetospheric plasmas.
- Instrument most commonly used are ESAs, Faraday cups and SSDs.
- ESAs and FCs are *energy/charge detectors*.
- Solid state detectors are total energy detectors, mainly used for detecting higher energy particles. *New* SSDs can measure particles from *a few keV* to several MeV and higher.
- We limit discussion of how ESAs work (See Wüest et al., 2007 for other types of detectors)

# Cluster and Wind Ion instruments



- A schematic diagram of a symmetric spherical “top hat” ESA (Carlson et al., 1987).
- 3D information obtained in one spin of the spacecraft.
- Cluster and Wind instrument (Lin et al., 1995; Réme et al., 1997).

- Concentric spheres have a mean radius  $R$ .  
electric field  $E$  applied between the plates.

- Particles travel in circular path will pass through the plates only if the *electric force just balances the centripetal force*,

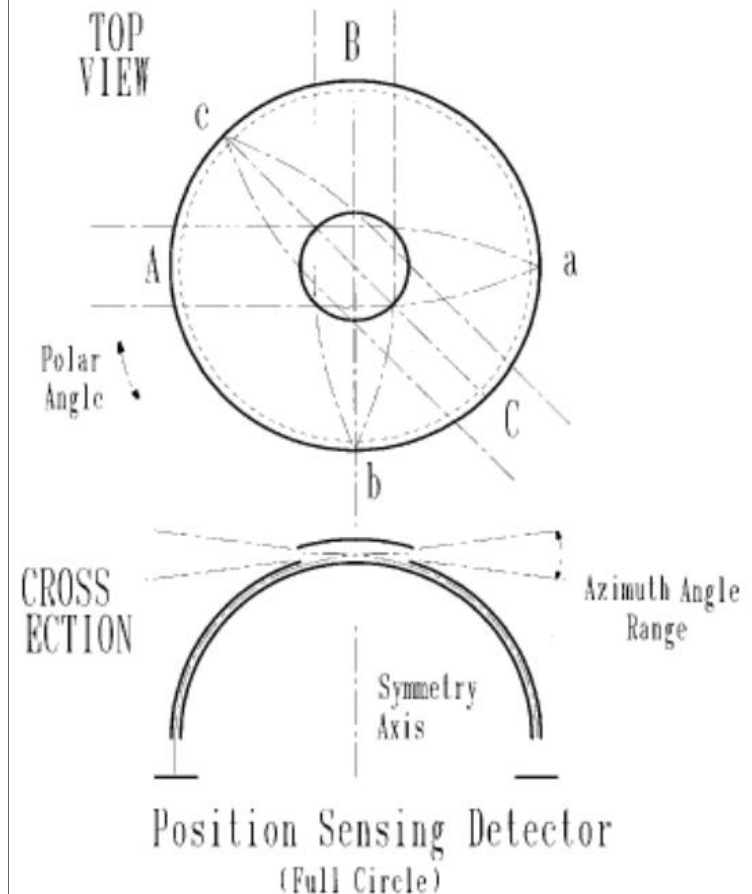
$$mv^2/R = qE$$

- Rewrite this equation,

$$mv^2/2q = ER/2$$

where *energy/charge* (left side) is related to instrument quantities (voltage & radius) on the right.

- ESAs measure *energy/charge* of the particle, regardless of the mass, charge or velocity.



*Important FACT:*

- Immerse ESA in plasma of average density  $n$  where the particles move with a mean velocity  $\langle \mathbf{v} \rangle$ . Total particle flux entering the aperture of a detector is  $n \langle \mathbf{v} \rangle$ , where

$$n(\mathbf{r}) = \int f(\mathbf{r}, \mathbf{v}) d^3v$$

$$\langle \mathbf{v} \rangle = \int \mathbf{v} f(\mathbf{r}, \mathbf{v}) d^3v$$

Here  $f(\mathbf{r}, \mathbf{v})$  is the distribution function of the particles.

- A detector measures the product  $n \langle \mathbf{v} \rangle$ , **not  $n$  or  $\langle \mathbf{v} \rangle$** .



- *A detector* counts particles. The total count is

$$C = n \langle v \rangle A \Delta t,$$

where  $A$  is the effective area of the entrance aperture and  $\Delta t$  is the accumulation time.

- Energy/charge (*E/q spectrum*) is obtained by measuring the particles over a small energy range  $\Delta E$  (Note  $E$  used for both electric field and energy).
- Define a *differential number of counts*:

$$C_i = n_i \langle v \rangle_i A \Delta E_i \Delta t$$

where  $C_i$  represents counts in the HV step  $i$  covering the narrow energy range  $\Delta E_i$ .

- Total *E/q spectrum* over the entire energy range obtained by varying the high voltage (HV) applied between the plates. The number of energy steps for a typical ESA is 16, 32 or 64.

- *Differential number flux*

$$F_N = C/g_v E \Delta t = C/g_E E \Delta t$$

Units: (cm<sup>-2</sup>-s<sup>-1</sup>-sr<sup>-1</sup>-eV<sup>-1</sup>), g<sub>E</sub> (cm<sup>2</sup>-sr), E in eV or keV.

- *Energy flux*

$$F_E = C/g_E \Delta t \quad \text{Units: ergs/cm}^2\text{-s}$$

- *Distribution function*

$$f(v) = C/\Delta t g_E v^4 \quad \text{Units: s}^3\text{-cm}^{-6}.$$

$F_N$ ,  $F_E$ , and  $f(\mathbf{r}, \mathbf{v})$  are *primary quantities* measured by instruments.

*Macroscopic quantities* are *computed* from measured quantities:

$$n(\mathbf{r}, t) = \int f(\mathbf{r}, \mathbf{v}, t) d^3v$$

$$\langle \mathbf{v} \rangle = \int \mathbf{v} f(\mathbf{r}, \mathbf{v}, t) d^3v$$

$$\langle v^2 \rangle = \int v^2 f(\mathbf{r}, \mathbf{v}, t) d^3v$$

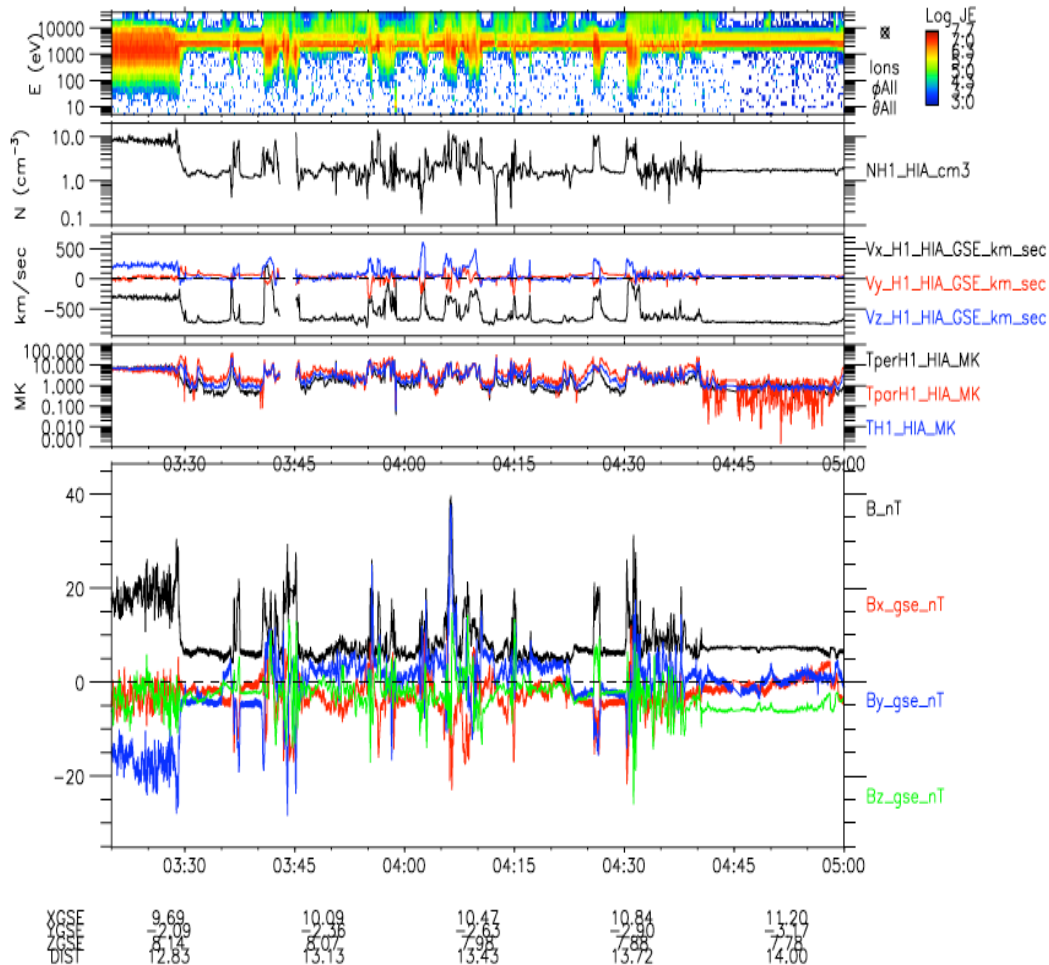
$$\mathbf{P} = m \int (\mathbf{v} - \langle \mathbf{v} \rangle)(\mathbf{v} - \langle \mathbf{v} \rangle) f(\mathbf{r}, \mathbf{v}, t) d^3v$$

$$\mathbf{P} = n k T$$

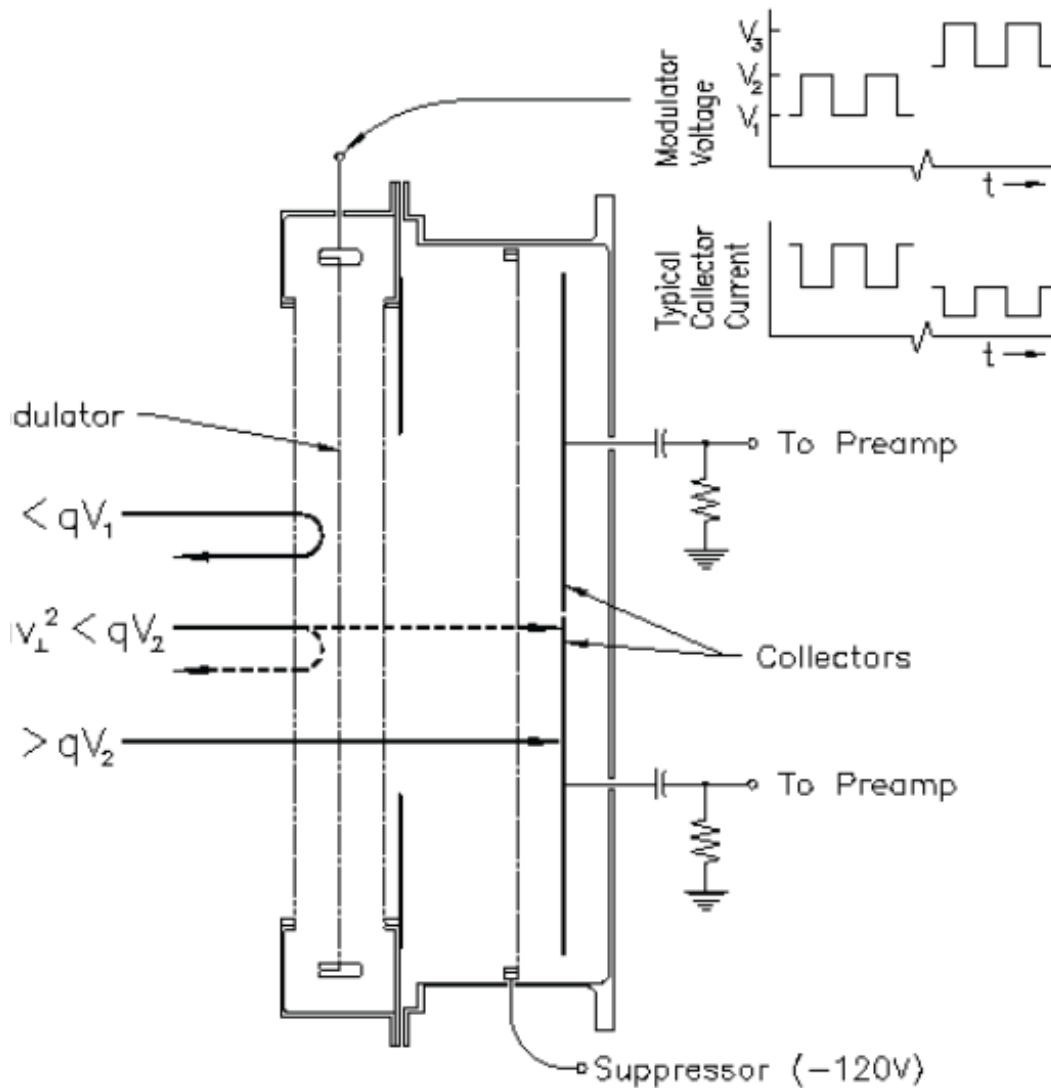
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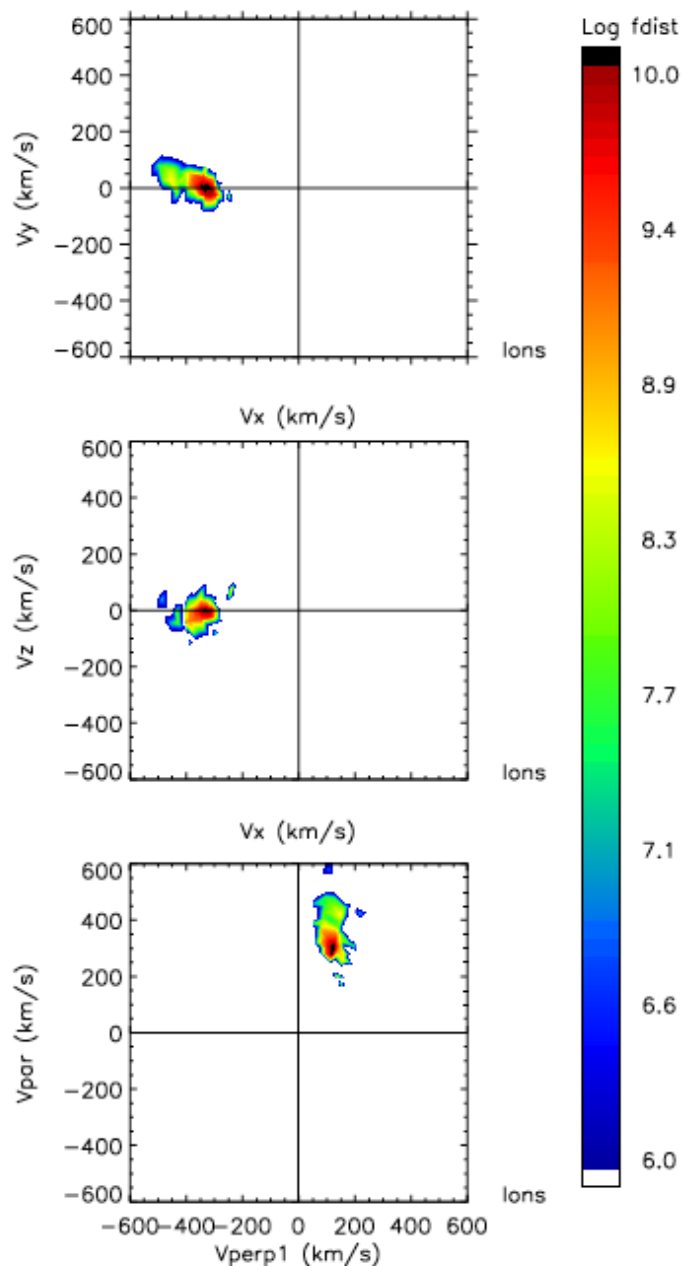
- Typical summary Plot from an ESA (Cluster)
- Energy flux (top) and computed Bulk parameters  $n$ ,  $\langle v \rangle$  and  $T$ .
- Magnetic field (bottom)



- Faraday cups measure particles from a few tens of eV to a few keV.

- Faraday Cups rugged, can operate for many years (Voyager)

$$I = eA \int \mathbf{v} f(\mathbf{v}) S(\mathbf{v}) d^3v$$



- ESAs do *not* measure  $v$ ,  $q$ , or  $m$ . However, SW data are plotted in velocity space, identifying  $H^+$  and  $He^{++}$  ion beams.
- How is information on  $v$ ,  $q$  and  $m$  of the different particles obtained?

- Energy per charge of  $H^+$  and  $He^{++}$  ions ( $\alpha$ 's) are

$$H^+ \quad (E/q)_+ = m_+ v_+^2 / 2q_+$$

$$He^{++} \quad (E/q)_\alpha = m_\alpha v_\alpha^2 / 2q_\alpha = m_+ v_\alpha^2 / q_+$$

where  $m_\alpha = 4m_+$  and  $q_\alpha = 2q_+$ .

- Interpretation of ESA data has assumed that

*“all particles are traveling at the same mean velocity in steady-state plasmas with a frozen in magnetic field”*

*Hundhausen, 1968*

- If  $H^+$  and  $He^{++}$  are *traveling together*, then  $v_+ = v_\alpha = V_{sw}$ .

For  $H^+$

$$(E/q)_+ = m_+ v_+^2 / 2q_+ = m_+ V_{sw}^2 / 2q_+$$

For  $He^{++}$

$$(E/q)_\alpha = m_+ v_\alpha^2 / q_+ = m_+ V_{sw}^2 / q_+$$

Hence,

$$(E/q)_\alpha = 2 (E/q)_+$$

- Thus, if we *assume* all particles are  $H^+$  in the velocity space, find a beam centered at  $V_{sw}$  and identify it as  $H^+$ . Another “ $H^+$ ” beam centered at  $(2)^{1/2} V_{sw}$  will be interpreted as  $He^{++}$  ions.

- A mass analyzer is needed to identify  $v$  and  $m/q$ .



*Basic Theories and Concepts* to interpret space plasma observations:

1. Coupled Lorentz-Maxwell equations (6N equations).
  2. Coupled Boltzmann-Maxwell equations (Use distribution function)
  3. Coupled Fluid-Maxwell equations (Use macroscopic variables)
- 1 and 2 *are equivalent* for *collisionless plasmas*. Theory is self-consistent and gives a complete picture of space plasma.
  - Lorentz-Maxwell approach avoided in the past because analytical solutions not possible.
  - Today, the coupled theory used more often because we have super computers to track the particles.
  - Most PIC simulations limited to 1 and 2D as computer capability still limited. However, computer capability is continually improving.
  - Simulation tools important for data analysis to help interpret complex features.
  - MHD fluid equations are conservation equations obtained from the *velocity moments* of the Boltzmann equation. They describe an approximate picture.

## Basic theory of space plasmas

The coupled Lorentz (Boltzmann) and Maxwell equations

$$\begin{aligned}\frac{d\mathbf{p}_k}{dt} &= q_k(\mathbf{E} + \mathbf{v}_k \times \mathbf{B}) \\ \frac{d\mathcal{E}_k}{dt} &= q_k \mathbf{v}_k \cdot \mathbf{E}\end{aligned}\quad (6)$$

$$\frac{\partial f(\mathbf{r}, \mathbf{v})}{\partial t} + \mathbf{v} \cdot \frac{\partial f(\mathbf{r}, \mathbf{v})}{\partial \mathbf{r}} + \frac{q}{m}(\mathbf{E} + \mathbf{v} \times \mathbf{B}) \cdot \frac{\partial f(\mathbf{r}, \mathbf{v})}{\partial \mathbf{v}} = 0 \quad (7)$$

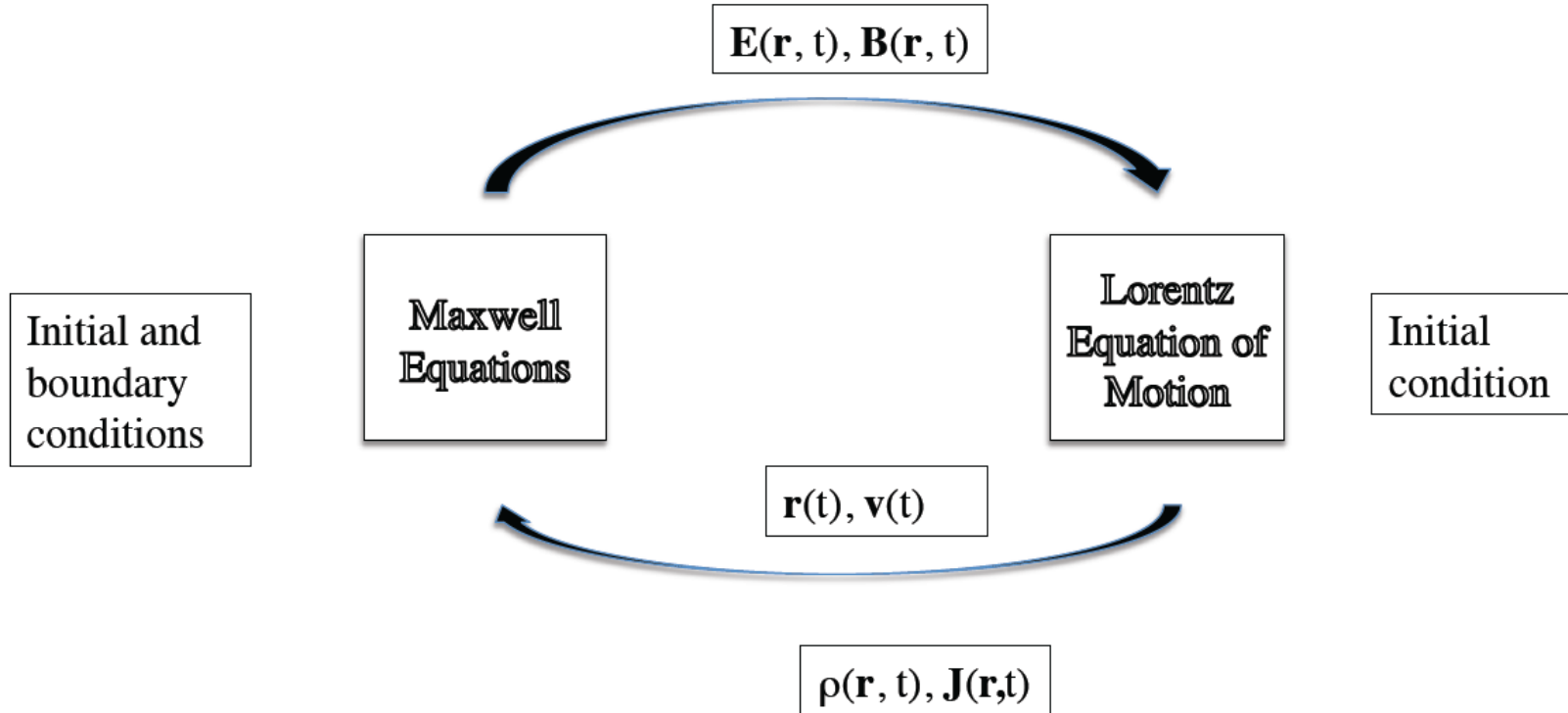
$$\nabla \cdot \mathbf{B} = 0 \quad (8)$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \partial \mathbf{D} / \partial t \quad (9)$$

$$\nabla \cdot \mathbf{D} = \rho_c \quad (10)$$

$$\nabla \times \mathbf{E} = -\partial \mathbf{B} / \partial t \quad (11)$$

## Self-consistent theory of space plasmas



- *Self-Consistency*: Particle motions produce the required electromagnetic fields that in turn are necessary to create the particle motions.

## MHD Equations:

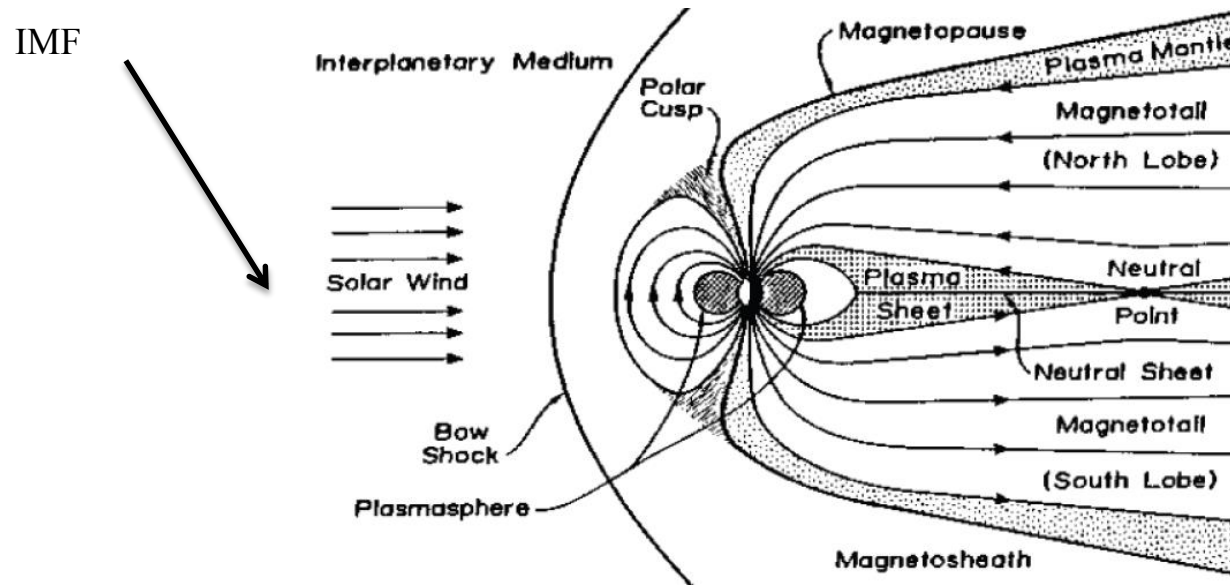
- The first three velocity moments yield *mass, momentum and energy conservation* equations.
- *Advantages*: Reduces the number of variables from  $6N$  to a few macroscopic variables:  
 $n, \mathbf{U}, T, \dots$

$$\frac{\partial n}{\partial t} + \nabla \cdot n\mathbf{U} = 0 \quad (1)$$

$$d\mathbf{U}/dt = -\nabla p + \mathbf{J}' \mathbf{B} \quad (2)$$

$$\frac{\partial}{\partial t} [nmU^2/2 + p/(g-1) + B^2/2m_0] + \nabla \cdot [nmU^2\mathbf{U}/2 + gp\mathbf{U}/g-1 + E\mathbf{x}B/m_0] = 0 \quad (3)$$

## *MHD Description of Solar Wind , IMF, bow shock, and Magnetosphere*



- SW *flows* out from the Sun.
- Solar magnetic field *transported* out *frozen* in the SW.
- SW is *supersonic*, hence a *shock wave* forms in front of Earth.
- Magnetosphere formed by the SW confining the geomagnetic field.
- A long tail produced by convecting “connected” IMF-geomagnetic field with the SW.

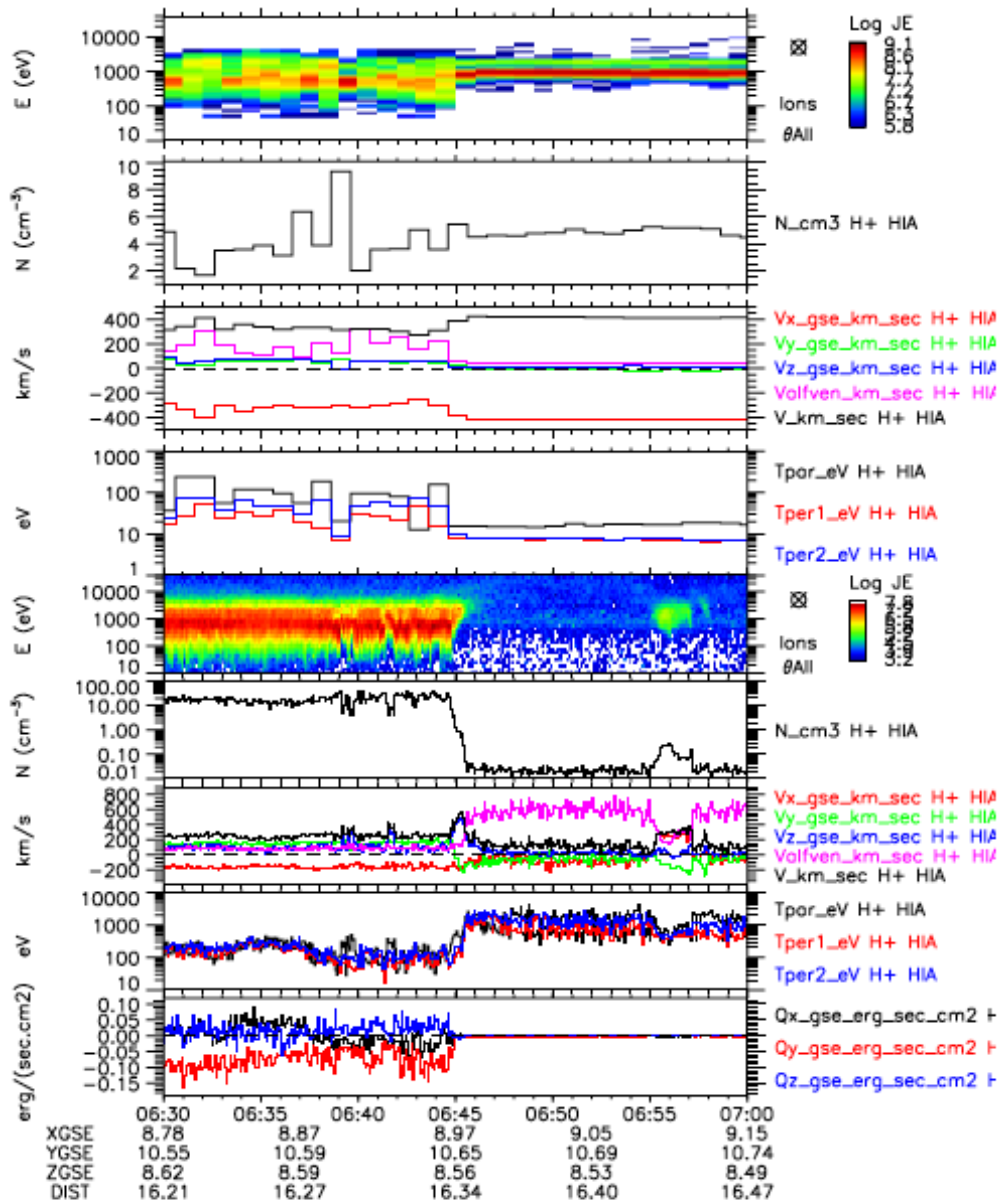
- MHD equations alone *not sufficient* to describe space plasma behavior self-consistently.
- There are always *more unknowns* than *number of equations*.
- For example, Particle flux conservation:  $\partial n / \partial t + \nabla \cdot n \mathbf{U} = 0$ ,  
Four unknowns (*n*, *U*), only three equations.
- Computing higher moments does not solve the problem. *New unknowns* are introduced.
- For a complete MHD description, one needs all velocity moments to solve the closure problem. *Not practical!*

- For a *finite number* of moment equations, MHD equations often supplemented by *Adiabatic equation of state* or *Ohm's law*.
- Adiabatic plasma: *No heat flux*, hence not consistent with many space plasma observations.
- Ohm's law. *No conductivity model exists for collisionless plasmas*.
- To remedy this problem, MHD treats space plasmas as fluid with *infinite conductivity* ( $\sigma = \infty$ ).
  - *Ideal fluids* conserve magnetic flux, leads to frozen-in-field dynamics: *No EMF* is generated.
- Approximation means you throw away information. You need to ask *what and how much physics is lost*.

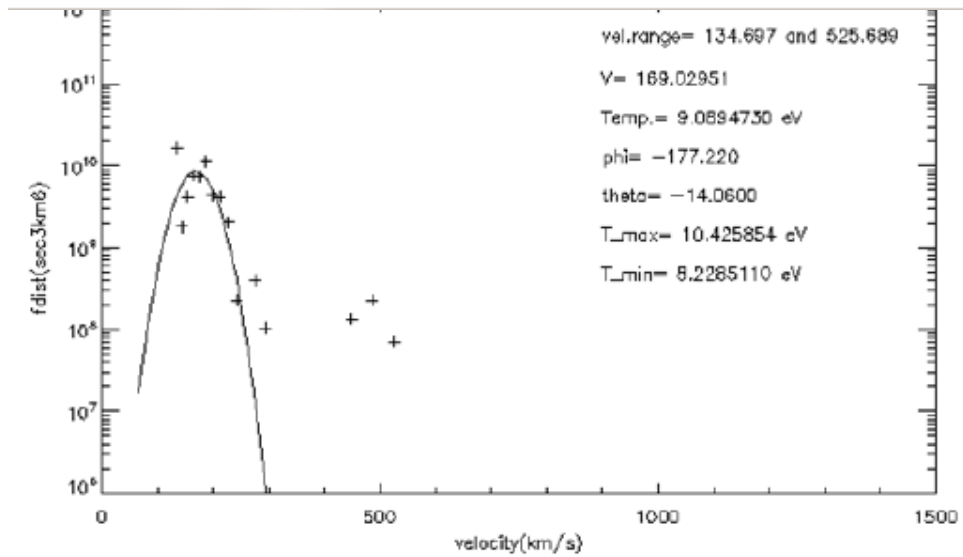
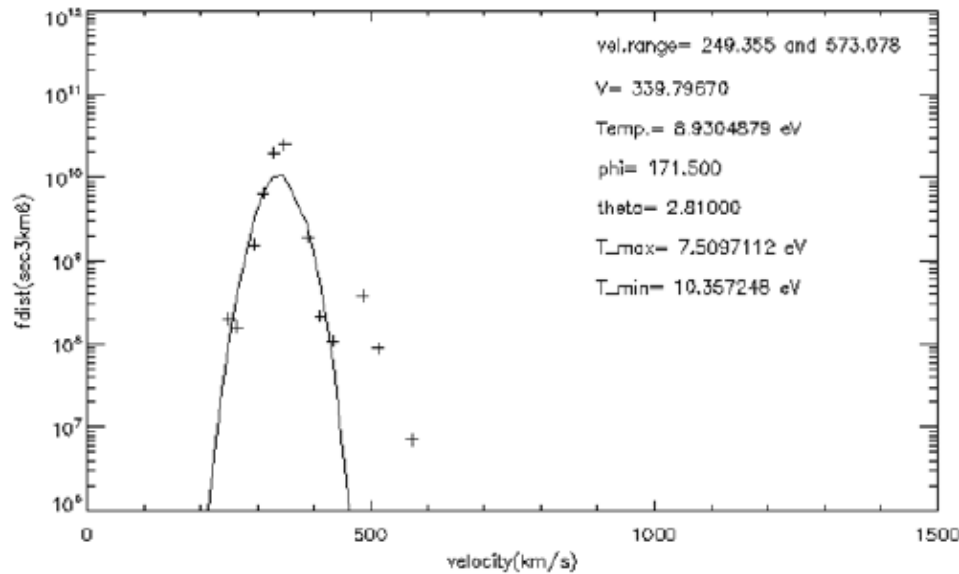
- Observations *not fully explained* by MHD theories and concepts.
- *Solar Wind: Heat flux carried by electrons.*
- *Bow shock*: Different from ordinary fluid shocks. Bow shock reflects up to 20% of incident SW back into the upstream region
- The remaining 80% transmitted across bow shock is not immediately thermalized.
- The bulk flow in the downstream of bow shock can often remain *super-Alfvenic*.



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- Bulk flow remains Super-Alfvénic in Magnetosheath.
- SW is not thermalized at the bow shock.



- SW  $H^+$  beam *slowed down* going across the bow shock but *not* thermalized.
- What shifts down the peak of the SW beam?

*Fundamental equation for Electric Field.*

From  $\nabla \times \mathbf{E} = -\partial \mathbf{B} / \partial t$  and  $\mathbf{B} = \nabla \times \mathbf{A}$ , obtain

$$\nabla \times (\mathbf{E} - \partial \mathbf{A} / \partial t) = 0. \text{ Let } \mathbf{E} = -\nabla \phi, \text{ then}$$

$$\mathbf{E} = -\nabla \phi - \partial \mathbf{A} / \partial t \quad (13)$$

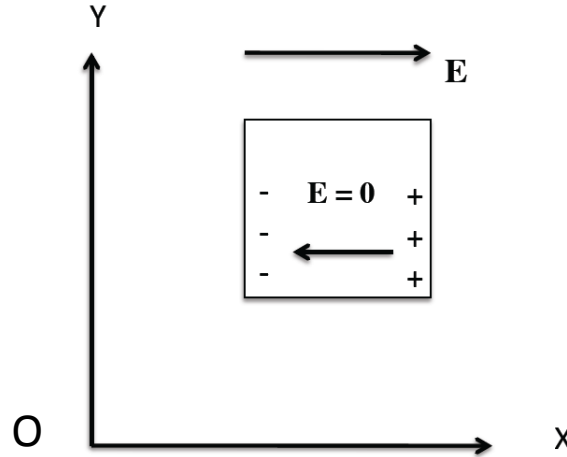
where

$$\phi(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{(\rho', t) d^3 r'}{|\mathbf{r} - \mathbf{r}'|} \quad (14)$$

and

$$\mathbf{A}(\mathbf{r}, t) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}', t) d^3 r'}{|\mathbf{r} - \mathbf{r}'|} \quad (15)$$

## Illustrate how particles in plasmas respond to electric force

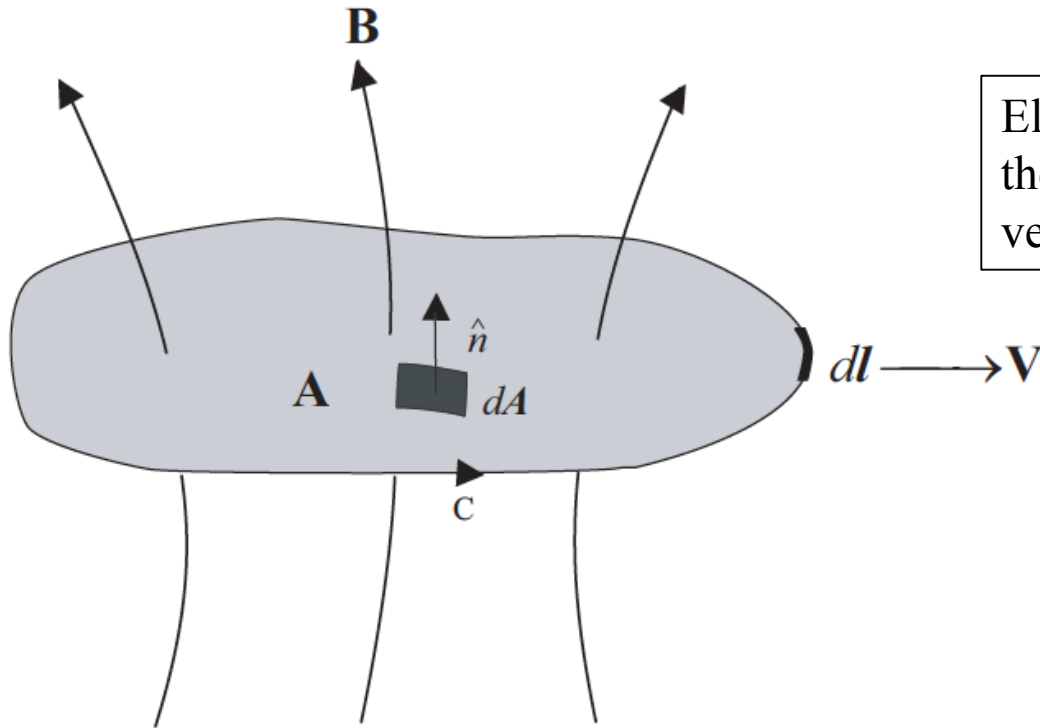


- Let an isolated plasma blob be uniform in space and charge neutral with equal number of protons(  $p^+$ ) and electrons ( $e^-$ ). The plasma blob is in equilibrium.
- Apply an  $E$ -field to a *stationary plasma blob*. No magnetic field,  $B=0$
- Inside the plasma blob, the force  $qE$  pushes electrons and ions in opposite directions. Produces an  $E$ -field opposing applied  $E$ .
- Motion stops when the *total force* on the particles *vanishes*.
- *$E = 0$  inside equilibrium plasmas*

$$\mathbf{E} = -\nabla\phi - \partial\mathbf{A}/\partial t$$

- The first term requires free charges in plasma. Plasmas have high *electrical conductivity*. No free charges accumulate so the first term disappears (*Caveat*: Free charges  $\rho$  can exist in *double layers*).
- *Inductive electric fields* responsible for the dynamics of space plasmas.

*Faraday's law*, one of the most important equations for space plasmas



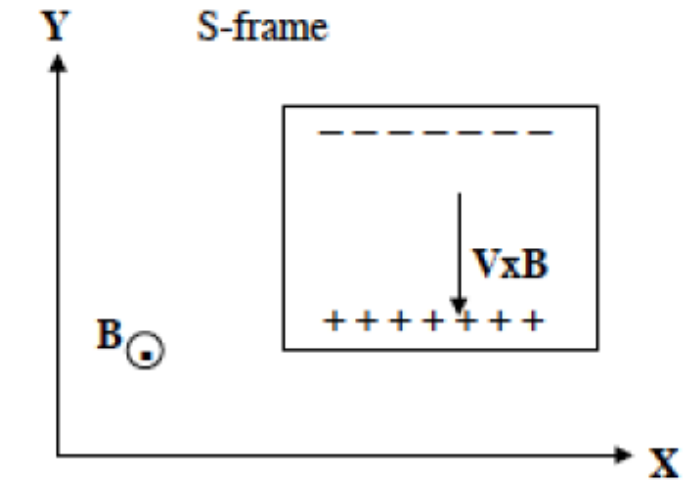
Electric field measured in the frame of the contour element  $dl$  moving with velocity  $V$  ( $S'$  frame).

$$EMF = -d\Phi/dt$$

$\Phi$  = magnetic flux enclosed by the contour C

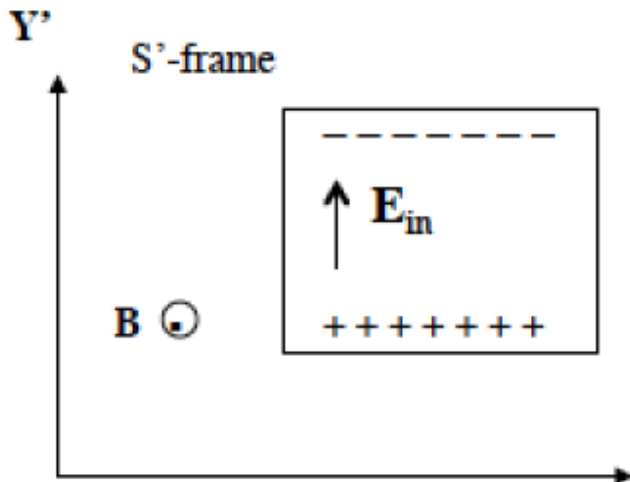
$$\oint (\mathbf{E} + \mathbf{V} \times \mathbf{B}) \cdot d\mathbf{l} = EMF$$

Moving Plasma Blob.  $\mathbf{B} \neq 0$ ,  $\mathbf{E} = 0$



$\mathbf{V}$   $\longrightarrow$

$\mathbf{V} = \mathbf{V}_X$   
 $\mathbf{F} = q\mathbf{V} \times \mathbf{B}$   
 Charge separation  
 Motion Until  $\mathbf{F} = q(\mathbf{E} + \mathbf{V} \times \mathbf{B}) = 0$



$\mathbf{V} = 0$   
 $\mathbf{F} = q\mathbf{V} \times \mathbf{B} = 0$   
 Charge separation still there  
 $\mathbf{F} = q\mathbf{E}_{in}$

*Moving Space plasmas.* Physics can be examined in stationary and moving frames. The quantities in different coordinate systems given by Lorentz transformation equations.

Lorentz transformation equations for  $\mathbf{E}$  and  $\mathbf{B}$  can be written vectorially as

$$\begin{aligned} \mathbf{E}'_{\parallel} &= \mathbf{E}_{\parallel} \\ \mathbf{E}'_{\perp} &= \gamma(\mathbf{E} + \mathbf{V} \times \mathbf{B})_{\perp} \\ \mathbf{B}'_{\parallel} &= \mathbf{B}_{\parallel} \\ \mathbf{B}'_{\perp} &= \gamma(\mathbf{B} - \frac{\mathbf{V}}{c^2} \times \mathbf{E})_{\perp} \end{aligned} \quad (12)$$

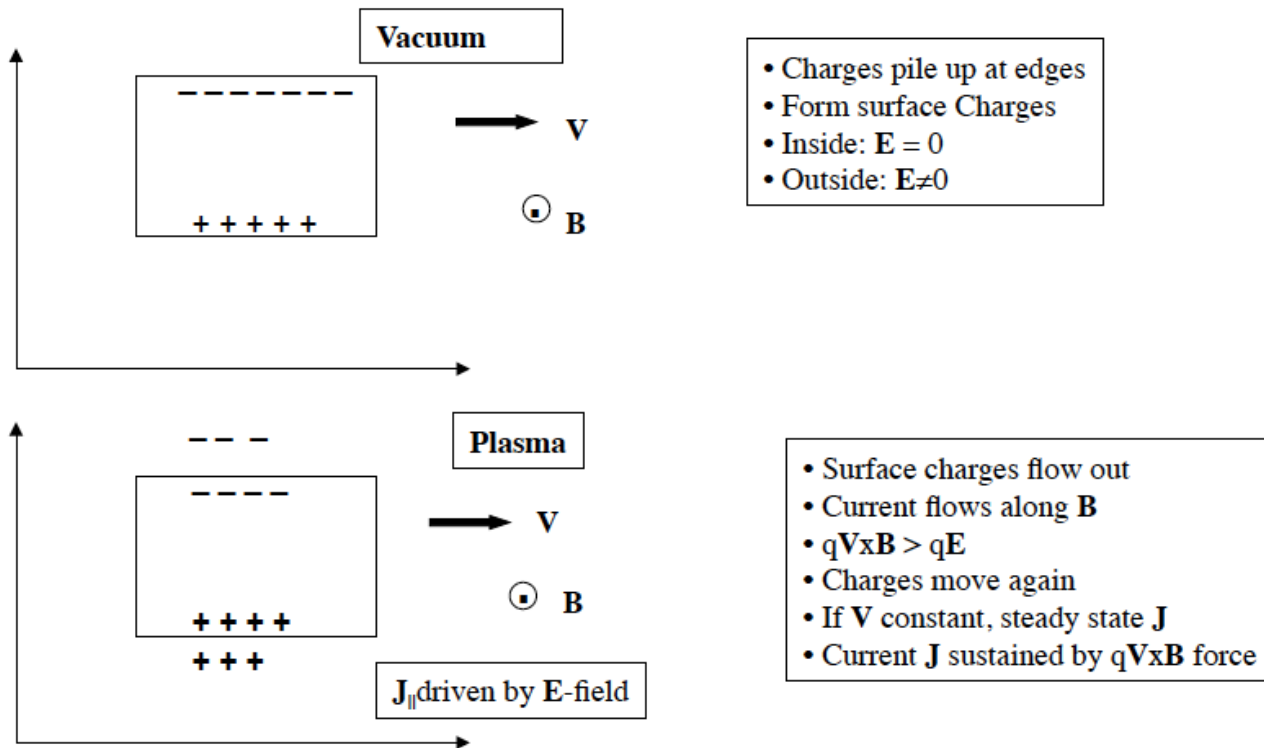
The subindices ( $\parallel$ ) and ( $\perp$ ) here refer to directions relative to  $\mathbf{V}$ , velocity of stationary frame relative to moving frame. For non-relativistic situation ( $\gamma = 1$ ) and to order  $V/c^2$ , the magnetic field is the same in the two frames but the electric field has a different expression,  $\mathbf{E}' = (\mathbf{E} + \mathbf{V} \times \mathbf{B})$ . Hence, when discussing electric fields, the reference frame must be stated.

$$E'_y = -V_x B_z, \quad B'_z = B_z \quad \text{and} \quad E'_y = E'_z = B'_x = B'_y = 0$$



# Motion of plasma blob surrounded by Vacuum (top) or by another plasma (bottom)

## Induced Current outside Blob



The End

# Vector Point Function

- The physical variables in Lorentz and Maxwell equation are *vector point functions*.

- Consider a point static charge  $q_k(\mathbf{r}) = q_k \delta(\mathbf{r} - \mathbf{r}_k)$ ,  
where  $\delta(\mathbf{r} - \mathbf{r}_k) = \delta(\mathbf{x} - \mathbf{x}_k) \delta(\mathbf{y} - \mathbf{y}_k) \delta(\mathbf{z} - \mathbf{z}_k)$ .

- *Electric field* produced by this charge given by Coulomb's law

$$\mathbf{E}(\mathbf{r}) = q_k (\mathbf{r} - \mathbf{r}_k) / |\mathbf{r} - \mathbf{r}_k|^3.$$

- Define Electric field at  $\mathbf{r}$  as *Force per unit charge*

$$\mathbf{E}(\mathbf{r}) = \lim_{q \rightarrow 0} \mathbf{F}_E / q$$

- $\mathbf{E}$  is parallel to  $\mathbf{F}_E$  and the charge  $q$  is accelerated in the direction  $\mathbf{E}$ . If there are many charges,  $\mathbf{E}(\mathbf{r}) = \sum q_k (\mathbf{r} - \mathbf{r}_k) / |\mathbf{r} - \mathbf{r}_k|^3$ .

- A set of point charges = charge density  $\rho(\mathbf{r}) = \sum q_k \delta(\mathbf{r} - \mathbf{r}_k)$ . Then

$$\mathbf{E}(\mathbf{r}) = \int d^3r \rho(\mathbf{r}) (\mathbf{r} - \mathbf{r}_k) / |\mathbf{r} - \mathbf{r}_k|^3.$$

- If a charge  $q$  is moving with velocity  $\mathbf{v}$ , there is now a *current*,  $q\mathbf{v}$ , which gives rise to a magnetic field  $\mathbf{B}$  at that point.
- In the presence of a magnetic field, a charge executes a circular motion due to the magnetic force,  $\mathbf{F}_B = q\mathbf{v} \times \mathbf{B}$ .
- Magnitude of  $\mathbf{F}_B$  depends on the magnitude and direction of  $\mathbf{v}$  and  $\mathbf{B}$  can be defined as *force per unit current*.

- Force is maximum when  $\mathbf{v}$  is perpendicular to  $\mathbf{B}$  and minimum when parallel to  $\mathbf{B}$ . The intensity of the magnetic field in terms of the maximum force  $|\mathbf{F}_B|_{\max}$  is

$$|\mathbf{B}| = \lim_{q\mathbf{v} \rightarrow 0} |\mathbf{F}_B|_{\max} / q\mathbf{v}$$

- The direction of  $\mathbf{B}$  is defined as the direction in which  $q$  would move when it experience no magnetic force.
- Continuous distribution of current, use Biot-Savat's law,
 
$$\mathbf{B}(\mathbf{r}) = \int d^3r' \mathbf{J}(\mathbf{r}') \frac{(\mathbf{r} - \mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|^3}.$$
- In the frame moving with the charge  $\mathbf{v}$ ,  $\mathbf{J}$  vanishes. But  $q$  is there and so is  $\mathbf{E}$ -field. One can thus look at  *$\mathbf{E}$  as primary quantity* and  *$\mathbf{B}$  is consequence of  $q$  in motion*.